

Supplementary Figure 1 | Nernst coefficient for LSCO cuprates.

Nernst coefficient *v* vs temperature for LSCO p = 0.15 (black), Nd-LSCO p = 0.15 (red), and Eu-LSCO p = 0.16 (green), plotted as v / T (figure adapted from ref. 1). The full lines are linear fits to the data and the arrows indicate the deviation from this linear behaviour at a temperature labeled T_v . The vertical dashed lines indicate the high temperature tetragonal (HTT) to orthorhombic (O) structural transition in LSCO (black), and the orthorhombic to low temperature tetragonal (LTT) transitions in Nd-LSCO (red) and Eu-LSCO (green). T_v is essentially the same for all three materials, independent of where the structural transitions occur.



Supplementary Figure 2 | Band structure calculations.

Band-structure calculations for Nd-LSCO in the LTT structure at ambient pressure (**a**) and in the HTT structure at P = 4.2 GPa (**b**). The lattice parameters for the calculations are taken directly from X-ray diffraction data: ref. 2 for (a) and ref. 3 for (b). A tight-binding fit to the band structure gives the following values for the parameters *t* and *t*' : *t* = 0.520 eV, *t*' = - 0.089, and *t*' / *t* = - 0.171, for LTT (P = 0) ; *t* = 0.539 eV, *t*' = - 0.082, and *t*' / *t* = - 0.152, for HTT (P = 4.2 GPa).



Supplementary Figure 3 | Hall coefficient vs doping in LSCO.

Hall coefficient $R_{\rm H}$ of LSCO in the T = 0 limit, as a function of doping p, for p = 0.18 and higher. Data are based on measurements made on crystals (red squares from ref. 4) and thin films (blue dots, from ref. 5). Lines are linear fits through the respective data. $R_{\rm H}$ is seen to decrease linearly as the doping increases away from $p_{\rm FS}$ (grey band). We use this linearity of $R_{\rm H}$ as a function of p to quantify the shift of $p_{\rm FS}$ with pressure in Nd-LSCO, in Figs. 4 and 5.



Supplementary Figure 4 | Effect of pressure on the Hall coefficient of Nd-LSCO at $p > p^*$.

a) Hall coefficient $R_{\rm H}$ of Nd-LSCO at p = 0.25 (above $p^* = 0.23$) as a function of temperature, measured at H = 16 T, for various applied pressures as indicated. b) Difference between two adjacent isobars in (a), for pairs as indicated. $R_{\rm H}(T)$ decreases rigidly, at a rate of - 6 % GPa⁻¹, relative to $R_{\rm H}(0) = 0.4$ mm³ C⁻¹.



Supplementary Figure 5 | Resistivity of LSCO under pressure.

Resistivity of LSCO as a function of temperature, in H = 45 T, for p = 0.143 (blue and dark blue) and p = 0.18 (red and burgundy), at pressures as indicated. Ambient pressure data for the same dopings, from ref. 6, are also shown, in grey for H = 0 and orange for H = 55 T.



Supplementary Figure 6 | Resistivity of Nd-LSCO at *p* = 0.15.

Resistivity of Nd-LSCO as a function of temperature, at p = 0.15 and in zero magnetic field, for three different pressures as indicated. Lines are linear fits at high temperature. The difference between data and fit is plotted in Fig. 2b.



Supplementary Figure 7 | Resistivity of Nd-LSCO at p = 0.23.

Electrical resistivity of Nd-LSCO at p = 0.23 as a function of temperature, under ambient pressure (blue) and 2.0 GPa (red), in H = 33 T. Grey data are ambient pressure and zero field data on the same sample. The black line is a linear fit to the 2.0 GPa data. ¹ Cyr-Choinière, O. *et al.* Pseudogap temperature *T** of cuprate superconductors from the Nernst effect. Preprint at http://lanl.arxiv.org/abs/1703.06927 (2017).

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